

An Integrated Bearing Parameter Identification Approach Using A Nonlinear Optimisation Scheme

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ABSTRACT

In system identification process, efficient global optimisation schemes play a vital role. In the present work, an effective optimisation-based identification methodology of bearing stiffness and damping coefficients using the bearing response data is proposed. The flexible rotor is initially analysed by finite element model with nonlinear bearing forces. The dynamic equations of rotor are solved to obtain the bearing responses in frequency domain at different operating speeds. The equivalent speed dependent parameters of stiffness and damping are then obtained by solving an error-based optimal formulation via modified particle-swarm optimisation scheme. The robustness of the approach is tested with added input noise. The approach is illustrated with the simulation results of rotor supported on full floating ring bearings and using an experimental data obtained from a rotor supported on two similar journal bearings. The error in identification is not exceeding five percent. The modified particle swarm optimisation is converging faster, and it is taking less computational time. The results are shown in the forms of graphs and tables.

Keywords: Bearing force coefficients; floating ring bearings; frequency response; hydrodynamic forces; nonlinear optimisation.

INTRODUCTION

Identification of bearing parameters is one of essential problems in rotor dynamics. The working speeds of certain industrial machinery are often very high, and, in such systems, even small vibrations persist for long time which leads to decrease in their life-span severely or even distorts the bearings fully. These high-speed rotors are often supported on various types of bearings and have unbalance and coupling forces leading to complex overall dynamics. The fluid film bearings are widely used in rotors due to their high damping and stiffness capabilities. Commonly linearised spring damper models are used at the bearing nodes for the system response and stability analysis of the rotor bearing system. Nevertheless, these parameters do not remain constant throughout the operating speed range and several approaches were proposed to identify such parameters. Changes in high speed rotor systems supported by nonlinear hydrodynamic forces of floating ring journal bearings were analysed thoroughly in literature [1-4]. Bearing parameters were obtained for plain and aerostatic journal bearings [5, 6]. Both linear and nonlinear bearing stiffness of journal bearing were estimated by perturbation technique using two-

dimensional Newton-Raphson iteration method on locating the journal equilibrium positions [7].

Methodology for prediction of sixteen dynamic coefficients for journal bearings in a rotor from experimental unbalance responses was presented earlier [8, 9]. Identification procedures of other bearing types are also found in various earlier works. A method with multi-frequency excitation for measurement of equivalent stiffness and damping of active magnetic bearing rotor was presented by Jiang et al. [10]. Based on the unbalance response [11, 12] and frequency characteristics [13], similar kinds of magnetic bearing parameter identification approaches were presented. Likewise for ball bearing systems, the parameters were identified with simulated and experimental data [14]. Likewise, linear and nonlinear bearing coefficients of oil-free bearings like gas-foil [15-19] and gas-film bearings [20, 21] have been also obtained. Response based identification methodologies for tilting pad journal bearings were proposed in literature [22-24]. Existing procedures for identification of bearing parameters were summarised in some review works [25, 26].

Hydrodynamic bearing forces are highly nonlinear and parametric in nature. Formulating an equivalent linear model is a highly involved procedure. For hydrodynamic journal bearings, a method of field estimation of parameters of stiffness and damping was illustrated [27] using measured responses at both shaft and housing locations. An experimental approach was proposed by Zhou et al. [28] to predict the parameters of stiffness and damping via the least-square minimisation under different operating conditions. A procedure to predict the rotor dynamic force coefficients of series bearing-supports was explored for impact and unbalance from the field measurements [29]. For estimation of bearing coefficients and unbalance from the measured responses, an optimisation based strategy was proposed by Kim et al. [30]. Kriging surrogate model together with differential evolution optimisation scheme was implemented [31, 32] to predict the bearing parameters. A modal parameter genetic time domain identification approach has been proposed [33] to know the characteristics of bearing using a multifrequency signal decomposition technology. Prediction of bearing parameter information chart is therefore very important task and a generalised methodology is therefore necessary to obtain the parameters conveniently.

Although many studies were found in literature, few works have only focused such estimation approach based on correlating the real time data with model-based outputs. In the present work, the speed dependent stiffness and damping parameters of the floating ring bearing system are obtained by minimising the average error in radial amplitudes between the measured and predicted values. Initially the rotor-bearing system is analysed by finite element model with nonlinear bearing forces and the unbalance response at bearing nodes is obtained at different speeds of operation. In the next stage, the rotor responses are generated by idealising the bearings as uncertain spring-damper system, whose parameters are to be correctly predicted. A modified particle swarm optimisation scheme is implemented to minimise the defined radial errors so as to achieve the correct set of stiffness and damping parameters. Robustness of methodology is tested with induction of noise to the measured data. The generalised approach is implemented further for prediction of bearing coefficients of an experimental rotor mounted on hydrodynamic journal bearings.

The paper is explained as follows: section-2 describes the mathematical modelling of rotor system using finite analysis, the expression of nonlinear bearing forces as well as the dynamic formulation of equivalent lumped parameter model of the rotor-bearing system. Section-3 presents the methodology employed and the proposed optimisation

scheme in brief. Section-4 gives the results and discussion of two cases considered in the analysis.

DYNAMIC MODEL OF A ROTOR BEARING SYSTEM

The dynamic model of the flexible rotor dynamic system is formulated using quasi-finite element analysis with lumped floating ring masses considered at the bearing locations. The rotor is treated as flexible component and disks are treated as rigid. Each individual node has 4 degrees of freedom (*DOF*) which includes two translations (v, w) and two bending slopes (φx , φy). The kinetic and potential energy expressions of the revolving shaft element can be expressed as below:

$$T_{el} = \int \frac{1}{2} \rho \left\{ A_s \left(\dot{v}^2 + \dot{w}^2 \right) + I_d \left(\dot{\phi}_x^2 + \dot{\phi}_y^2 \right) + I_p \left[\omega^2 + \omega \left(\dot{\phi}_x \phi_y - \dot{\phi}_y \phi_x \right) \right] \right\} ds$$
(1)

$$U_{el} = \int \frac{1}{2} \left\{ EI\left(\phi_{x}^{'2} + \phi_{y}^{'2}\right) + kGA\left[\left(\phi_{y} - v_{y}^{'}\right)^{2}\right] \right\} ds$$

$$\tag{2}$$

Each disk kinetic energy can be expressed as Eq. (3).

$$T_{d} = \frac{1}{2} m_{d} (\dot{v}^{2} + \dot{w}^{2}) + \frac{1}{2} J_{d} (\dot{\phi}_{x}^{2} + \phi_{y}^{2}) + \frac{1}{2} J_{p} \left[\omega^{2} + \omega \left(\phi_{x} \phi_{y} - \phi_{y} \phi_{x} \right) \right]$$
(3)

The virtual work done by unbalance forces at the disks can be expressed as Eq. (4).

$$\delta W = m_{d1} r_1 \omega^2 (v_1 \cos \omega t + w_1 \sin \omega t) + m_{d2} r_2 \omega^2 (v_2 \cos \omega t + w \sin \omega t)$$
(4)

Using Hamilton principle,

$$\delta \int (T - V + W) dt = 0 \tag{5}$$

with δ denote the variational symbol, the motion of equation of the rotor alone is written as:

$$M\ddot{u} + [C + \Omega G]\dot{u} + Ku = F_{un} + F_i$$
(6)

Here, *u* denote the displacement vector of size $Q \times I$, F_{un} represents both unbalance and gravity force vector at the disks. F_i is the inner oil film force. *M*, *C*, *G* and *K* denote respectively the assembled system mass, system damping, gyroscopic and stiffness matrices of size $Q \times Q$ The motion of equation of ring is expressed with the help of inner and outer oil film forces, feed pressure of lubricant and floating ring dead weight. The final motion of equation for the floating ring is written as:

$$M_R \ddot{u}_R = F_o - F_i \tag{7}$$

where u_R denote the displacement vector of the floating rings at two bearings with the size 4×1 , M_R = diag(m_R , m_R , m_R , m_R , m_R) is the mass matrix of the rings with size of 4×4 . The outer and inner lubricant film forces of floating ring bearing can be expressed as Fo and Fi. The final assembled equation of motion after combining the rotor and bearing system of equations can be written as follows:

(8)

Mü+Cü+Kũ=F

where $\tilde{u}=[u \ u_R]^T$ is displacement vector matrix of size $(Q+4)\times 1$ and $\tilde{F}=[F_{un}+F_i \ F_0-F_i]^T=\tilde{F}_{un}+\tilde{F}_b$ is bearing forces and unbalance resultant force vector of size $(Q+4)\times 1$. while $\tilde{M}=\begin{bmatrix}M & 0\\0 & M_R\end{bmatrix}$, $\tilde{C}=\begin{bmatrix}C+\Omega G & 0\\0 & 0\end{bmatrix}$, $\tilde{K}=\begin{bmatrix}K & 0\\0 & 0\end{bmatrix}$ are effective square matrices of size $(Q+4)\times(Q+4)$.

Forces of Bearing

A floating-ring bearing consists of an annular ring placed between the sleeve and journal and there is a thin oil film in between them. The bearing mid-plane consists of the circumferential feed grooves and the lubricant is fed from the journal position to sleeve through bunch of feed holes located on the ring. The coordinate system considered is shown in Figure 1.



Figure 1. Geometry of the floating-ring bearing.

The hydrodynamic oil film forces can be obtained from the 2-D Reynolds equation with the use of the lubricant pressure distribution. Reynolds equations for both inner and outer lubricant films can be expressed as:

$$\frac{1}{R_{jr}^2} \frac{\partial}{\partial \varphi_{in}} \left(\frac{h_{in}^3}{12\mu_{in}} \frac{\partial p_{in}}{\partial \varphi_{in}} \right) + \frac{\partial}{\partial z_{in}} \left(\frac{h_{in}^3}{12\mu_{in}} \frac{\partial p_{in}}{\partial z_{in}} \right) = \frac{\Omega_{jr} + \Omega_m}{2} \frac{\partial h_{in}}{\partial \varphi_{in}} + \frac{\partial h_{in}}{\partial t}$$
(9)

$$\frac{1}{R_{ou}^2} \frac{\partial}{\partial \phi_{ou}} \left(\frac{h_{ou}^3}{12\mu_{ou}} \frac{\partial p_{ou}}{\partial z_{ou}} \right) + \frac{\partial}{\partial z_{ou}} \left(\frac{h_{ou}^3}{12\mu_{ou}} \frac{\partial p_{ou}}{\partial z_{ou}} \right) = \frac{\Omega_m}{2} \frac{\partial h_{ou}}{\partial \phi_{ou}} + \frac{\partial h_{ou}}{\partial t}$$
(10)

where p and μ represents the lubricant film pressure and viscosity; the subscripts in and ou denote the parameters of inner and outer oil films, respectively. Journal and floating ring bearing parameters are differentiate with the subscripts jr and rn. The journal and floating ring outer radius represented as R_{jr} and R_{ou}. The angular coordinated for inner and outer lubricant film is represented as ϕ . The axial coordinates of the inner and outer films are denoted by z_{in} and z_{ou} respectively. The oil film thicknesses h_{in}, h_{ou} and the squeezed terms $\partial h_{in}/\partial t$, $\partial h_{ou}/\partial t$ can be expressed by considering short length bearing theory as follows:

$$h_{in}(\varphi_{in},t) = C_{in} - x_{jr} \cos \varphi_{in} - y_{ir} \sin \varphi_{in}$$
(11)

$$h_{ou}(\varphi_{ou},t) = C_{ou} - X_{rn} \cos \varphi_{ou} - Y_{rn} \sin \varphi_{oo}$$
(12)

$$\frac{\partial h_{in}}{\partial t} = -(\dot{x}_{jr}\cos\varphi_{in} - \dot{y}_{jr}\sin\varphi_{in})$$
(13)

$$\frac{\partial h_{ou}}{\partial t} = -(\dot{X}_{rn} \cos \varphi_{ou} - \dot{Y}_{rn} \sin \varphi_{o})$$
(14)

$$\begin{cases} x_{jr} = X_{jr} - X_{rn} \\ y_{jr} = Y_{jr} - Y_{rn} \end{cases}$$
(15)

$$\begin{cases} \dot{x}_{jr} = \dot{X}_{jr} - \dot{X}_{rn} \\ \dot{y}_{jr} = \dot{Y}_{jr} - \dot{Y}_{rn} \end{cases}$$
(16)

where (X_{jr}, Y_{jr}) are the journal center O_{jr} displacement vector in the fixed reference frame. Also, C_{in} and C_{ou} represent the inner and outer oil film static clearances. The absolute displacement and velocity components of the floating ring center O_{rn} , are (X_{rn}, Y_{rn}) and $(\dot{X}_{rn}, \dot{Y}_{rn})$ while the absolute velocities of centers O_{jr} are denoted as $(\dot{X}_{jr}, \dot{Y}_{jr})$. The lower case letters (x_{jr}, y_{jr}) , $(\dot{x}_{jr}, \dot{y}_{jr})$ denote the displacement and velocity components of O_{jr} relative to O_{rn} . The expressions for oil film forces with short bearing approximation can be written as [34]:

$$\begin{cases} F_{ix} \\ F_{iy} \end{cases} = \mu_{in} (\Omega_{jr} + \Omega_{rn}) R_{jr} L_{in} \left(\frac{R_{jr}}{C_{in}}\right)^2 \left(\frac{L_{in}}{2R_{jr}}\right)^2 \begin{cases} f_{ix} \\ f_{iy} \end{cases}$$
(17)

where f_{ix} , f_{iy} and f_{ox} , f_{oy} respectively denote the component forces on the journal and bearing. Detailed expressions of these forces are given in the Appendix-A.

Methodology and Optimisation procedure

The bearing force components are expressed in terms of displacements and velocities in bearing coordinates as:

$$\begin{cases} \tilde{F}_{Bx} \\ \tilde{F}_{By} \end{cases} = \begin{bmatrix} c_{xx} & c_{xy} \\ c_{yx} & c_{yy} \end{bmatrix} \begin{pmatrix} \dot{x}_B \\ \dot{y}_B \end{pmatrix} + \begin{bmatrix} k_{xx} & k_{xy} \\ k_{yx} & k_{yy} \end{bmatrix} \begin{pmatrix} x_B \\ y_B \end{pmatrix}$$
(19)

where the terms c and k represent the unknown direct and cross-coupled damping and stiffness bearing force coefficients and the suffix B denotes the bearing support location. By substituting these bearing forces in equation (8) and converting into frequency domain, the system of equations can be written as:

$$\begin{cases} F_{x} \\ F_{y} \end{cases} - \omega^{2} \widetilde{M} \begin{cases} \& X \\ \& Y \end{cases} = \begin{bmatrix} H_{xx} & H_{xy} \\ H_{yx} & H_{yy} \end{bmatrix} \begin{cases} \& X \\ \& Y \end{cases}$$
 (20)

Here, $H_{ij}(\omega) = k_{ij}(\omega) + i\omega c_{ij}(\omega)$ is impedance function, (F_x, F_y) and (X,Y) are the discrete Fourier transforms of external forces and displacements respectively. With the knowledge of the component displacements in bending directions at two bearing nodes, it is possible to compute the eight direct coefficients corresponding to each of the two bearings. In order to obtain the correct set of parameters, an error function defined in terms of amplitudes of radial displacements at the bearing nodes is considered at every operating speed. Mathematically, the optimisation problem is defined as:

Minimise
$$E = (r_{1(FE)} - r_{1(L)})^2 + (r_{2(FE)} - r_{2(L)})^2$$
 (21)

where $r_{n(FE)} = \sqrt{X_{nFE}^2 + Y_{nFE}^2}$ is the amplitude at bearing-n in radial-direction from actual system, while $r_{n(L)} = \sqrt{X_{nL}^2 + Y_{nL}^2}$ is the radial displacement amplitude at bearing-n from linear bearing force model. This is an implicit function of bearing coefficients, whose upper and lower bounds are considered as: $k_i \in [k_{min}, k_{max}]$ and $c_i \in [c_{min}, c_{max}]$ with i=x,y. The minimisation problem is highly nonlinear function of the stiffness and damping parameters and requires a robust optimisation algorithm.

Particle Swarm Optimisation (PSO)

Particle swarm optimisation scheme is one of the robust metaheuristic optimisation schemes based on the behavior of flocking bird/fish in food search [35]. Initially with the unspecified set of solutions, the system starts and look for optimum solution by updating the generations. Particle is described as each candidate solution and set of particles is known as swarm. In a cooperative manner they move in n-dimensional search space. The variable which performs the Swarm movement is called velocity. This velocity is influenced by social and local factors. P_{best} is the best position identified so far in the search space and G_{best} is the best position identified by the swarm. In search space R^k (with k initial sets) for each particle, the objective function value is calculated. If S(n) and Vel(n) are each particle position and velocity at nth iteration, there are three major factors which influence the movements of each particle. Those are (i) own direction search of particle (ii) particle best position itself (iii) whole swarm best position. The position and velocity of every particle after iteration number n is updated using the following equation.

$$Vel(n+1) = u \times Vel(n) + c_1 r_1 (P_{best} - S(n)) + c_2 r_2 (G_{best} - S(n))$$
(22)

$$S(n+1)=S(n)+Vel(n+1)$$
 (23)

where u is called inertia factor of the particle, which often reduced in every cycle. Acceleration coefficients c_1 and c_2 describe the private (cognitive) and global (social) behavior of the system. Also, the random numbers r_1 and r_2 are considered in-between 0 and 1. Till the current iteration P_{best} is the particle best position. while G_{best} is the best position of the group until current iteration. The algorithm converges to a best swarm by selecting correct values of u, c_1 and c_2 . Further, the velocity of the particle is bounded

between minimum and maximum values. In order to improve accurate and faster convergence, a modification is employed.

Modified Particle Swarm Optimisation (MPSO)

In order to attain faster convergence without loss of accuracy, several modifications were suggested. To improve the population diversity and PSO's performance, mutation is a powerful strategy [36]. Here, a correction to the updated vector S(n+1) in every cycle is introduced. This approach evaluates a mutation vector created from randomly selected three swarms (vectors, Y_{B1} , Y_{B2} and Y_{B3}) in that generation. Mathematically, mutation vector (Y_{mut}) is expressed as [37]:

$$Y_{mut} = Y_{B1} + \gamma^* (Y_{B2} - Y_{B3})$$
(24)

where $\gamma * \in [0.9, 1]$ is the mutation constant The resultant population is modified using this mutation vector according to the following rule.

$$Y_{\text{new},i} = \begin{cases} Y_{\text{mut},i}, & \text{if}(\text{rand} < p_{\text{Cross}}) \\ Y_{i}, & \text{otherwise} \end{cases} \text{ for } i=1,2,...k.$$
(25)

Here, k is the number of points in the population (swarm size). The p_{cross} is crossover probability selected in the range of 0.1 to 0.9.

RESULTS AND DISCUSSION

The error between the actual response from the system and idealised response from the linear model is minimised by proper adjustment of the stiffness and damping parameters of the linear model. In order to illustrate the approach, three cases are considered. In the first case, the idealised system and actual system are both linear models, while in the second case actual system is nonlinear and in third case an experimental data is used as actual response. In all the simulated cases, the rotor is analysed using finite element model with Timoshenko beam elements as shown in Figure 2. [1] The two disc masses are mounted at nodes 1 and 9, while the nodes 3 and 7 denote the bearings. Dimensional data of the rotor-bearing system is depicted in Table 1.



Figure 2. FE model considered for analysis.

Properties	Value
Shaft material density, ρ (kg/m ³)	7800
Left disk mass, M _{DL} (kg)	1.4
Right disk mass, M _{DR} (kg)	1
Left disk dimeter moment of inertia, I _{DL} (kgm ²)	6.3×10 ⁻⁴
Right disk diameter moment of inertia, I _{DR} (kgm ²)	4.5×10^{-4}
Left disk polar moment of inertia, J _{DL} (kgm ²)	1.26×10 ⁻⁵
Right disk polar moment of inertia, J _{DR} (kgm ²)	9×10 ⁻⁴
Rotor diameter, D _{rotor} (m)	0.02
Rotor Length (m)	0.4
Young's modulus, E (GPa)	200
Radius of bearing (m)	0.01
Length of bearing (m)	0.01
Radial clearance of bearing (microns)	200
Viscosity of oil film (Pa-s)	288×10 ⁻⁴
Eccentricity (m)	1×10 ⁻⁶

Table 1. Data and material properties of rotor-system [3].

The resultant equations are solved by fourth-order Runge-Kutta method with all zero initial conditions. All simulations are carried-out on an Intel core-i7 processor with 3.40 GHz speed and 4 GB RAM. All the codes are implemented in MATLAB R2015a. Figure 3 show the time domain response obtained with reference [38] stiffness and damping values.



Figure 3. Time domain response at left bearing at rotor speed of 5000 rpm

Test Case-1: Actual and Ideal Systems, Both are Linear

Initially the vibration response of the rotor from known linear bearing parameters is generated. This simulated data is used as a baseline for testing the accuracy of the proposed methodology. The PSO parameters are set as: $c_1=2.1$, $c_2=2.1$ and $u = u_{\text{max}} - j\left(\frac{u_{\text{max}} - u_{\text{min}}}{j_{\text{max}}}\right)$, where u_{max} is maximum weight u_{min} is minimum weight, j is

iteration number and j_{max} is maximum iterations. In present work u_{max}=0.9 and u_{min} =0.4 is employed [35]. The variable bounds are taken as: $k_{min}=1\times10^5$ N/m, $k_{max}=50\times10^5$ N/m, $c_{min}=1\times10^2$ N-s/m $c_{max}=50\times10^2$ N-s/m. Figure 4 shows the error as a function of swarm size at 5000 rpm. It is seen that the error is minimum at an optimal swarm size of 30. Table 2 shows the correct and predicted bearing parameters from the linear rotor model at a speed of 5000 rpm. The percentage error is found to be less.



Figure 4. Effect of swarm size in MPSO on final converged error.

Parameter	Reference	Without	With 5%	With 10%
	Values [38]	noise	noise	noise
k _{x1} (N/m)	2.50×10^{5}	2.532×10^{5}	2.554×10^{5}	2.591×10^{5}
k _{y1} (N/m)	2.75×10^{5}	2.763×10^{5}	2.799×10^{5}	2.825×10^{5}
k _{x2} (N/m)	2.75×10^{5}	2.778×10^{5}	2.785×10^{5}	2.798×10^{5}
k _{y2} (N/m)	2.82×10^{5}	2.832×10^{5}	2.892×10^{5}	2.934×10^{5}
c_{x1} (Ns/m)	300	303.24	303.94	304.64
c_{y1} (Ns/m)	399	400.91	401.15	402.27
c_{x2} (Ns/m)	315	316.98	317.84	319.58
c_{v2} (Ns/m)	300	302.54	302.81	303.47

Table 2. Assumed and predicted Bearing parameters for linear rotor model at 5000 rpm.

Figure 5 shows the convergence trend of fitness function using PSO and MPSO approaches. MPSO algorithm converges faster than PSO with same other parameters. In a similar way, the stiffness and damping coefficients are estimated with close accuracy at different rotor speeds.



Figure 5. Fitness function convergence in PSO and MPSO.

Test Case-2: System with Nonlinear Bearing Forces

Here simulated frequency domain data is obtained from FE model of rotor supported over floating-ring bearings. Table 3 shows the parameters of bearing employed in simulation. The time responses and frequency spectra at the bearing nodes in two directions are obtained in different speeds. Figure 6 shows the bearing node frequency spectra at a speed of 5000 rpm. There are two critical frequencies (83.62Hz and 160Hz) as observed from FFT spectra. Furthermore, the subharmonic resonances resulting from hydrodynamic bearing forces are relatively small at this speed of operation [39].

Table 3. Parameters of bearing considered [3].

Parameter	Value
Outer bearing clearance, $C_2(m)$	8×10 ⁻⁵
Inner bearing clearance, $C_1(m)$	2×10 ⁻⁵
Floating ring mass, mr (kg)	0.02
Inner oil film viscosity, μ_i (Pa-s)	0.006
Outer oil film viscosity, μ_0 (Pa-s)	0.012



Figure 6. Frequency spectra at 5000 rpm at (a) left and; (b) right bearing.

In order to identify the bearing coefficients these response data are given as inputs to the optimisation program at a speed of 5000 rpm. The modified particle swarm optimisation scheme is implemented with a swarm size of 30. In order to test the effectiveness, the minimisation of error function is carried out using two other metaheuristic algorithms namely genetic algorithms (GA) [40] and simulated annealing (SA) [41]. For GA, population sizes of 30 along with mutation and crossover functions respectively as constraint dependent and scattered type were selected. In SA also the function tolerance is considered as 10⁻⁶. The annealing function used is fast annealing type and the reannealing interval is considered to be 100. Figure 7 shows the convergence trends obtained from GA and SA.



Figure 7. Convergence trend from (a) GA (b) SA.

The error function value obtained along with the time taken from the two optimisation approaches are depicted in the Table 4 along with MPSO. It is seen that the MPSO achieves relatively a minimum error. The time taken is also small and it also permits minimum number of parameters to alter. The optimum stiffness and damping parameters obtained from the three techniques are found to be close to each other. Table 5 shows the identified stiffness and damping parameters of the bearings at different speeds of the rotor using MPSO. Using these identified linear coefficients, the response amplitudes at left bearing node are plotted as shown in the Figure 8. The amplitudes from input simulated data with nonlinear bearing force model are also depicted for comparison.

Sl.no	I	Error	Time (s)
1	MPSO	5.49e-17	140
2	GA	1.083e-14	155

185

1.082e-11

3

SA

Table 4. Error obtained from the three optimisation techniques.

The amplitudes are matching at lower speeds of operation in both the directions. Due to minimisation of the radial error, the identified amplitudes are close in Y-direction compared to X-direction beyond 5000 rpm. This may be due to the amplitudes in the Y-direction are of two orders of more magnitude those in X-direction, resulting in a Y-direction fitting for error function. This is due to nonlinear objective error function of radial amplitudes accounted in this work, instead of errors in individual directions. More unified non-dimensional amplitudes would lead to improved accuracy.



Table 5. Identified stiffness and damping parameters of the bearings at different speeds.

Figure 8. Comparison of vibration amplitudes at left bearing in X and Y-direction.

Test Case-3: Parameter Identification from Experimental Data

In order to further test the methodology, an experimental analysis is conducted on a dual disk rotor mounted over two oil-film bearings. Figure 9 shows the schematic of an inhouse fabricated rotor bearing system operated with a single-phase AC motor. Two accelerometers (PG 109 Mo, 1-10 kHz) are mounted on each bearing in lateral directions to measure the rotor displacements. Both the input and output signals are recorded with a four-channel digital storage oscilloscope (Tektronix, model- DPO 43034) for recording the signals. A similar set of observations are noted at the other bearing. Table 6 shows the geometric and material parameters of the rotor system. Two similar hydrodynamic journal bearings (SKF UCP 204, 12.8 KN) are employed for supporting the rotor.



Figure 9. Experimental rotor model employed.

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Properties	Value
Shaft material density (kg/m ³)	7800
Left disk mass, M _{DL} (kg)	1.4
Right disc mass, M _{DR} (kg)	1
Rotor diameter, $D_{sh}(m)$	0.016
Rotor length (m)	0.48
Young's modulus, E (GPa)	200
Distance between the bearings (m)	0.22
Distance from Disc1 to left bearing	0.09
Distance from Disc2 to right bearing	0.09

Table 6. Experimental system data.

Figure 10 shows the time history and FFT plots obtained from the experimental analysis at the left bearing node in two directions at rotor speed of 900 rpm. Likewise, the time history and FFT plots at the right bearing in two directions are also recorded. Unlike in simulation results, the operating speed is limited by the motor employed in the test setup which is up to 900 rpm. However, the identification approach can be tested at any speed of rotor operation.



Figure 10. Time history and FFT plots at the left bearing.

Using these frequency spectra, the amplitudes at the two bearings in X and Y directions are noted down. With proposed methodology of identification, stiffness and damping parameters are estimated using MPSO algorithm. Figure 11 shows the estimated linear stiffness and damping parameters at the two bearing nodes corresponding to 900

rpm. Figure 12 shows the frequency responses obtained from the measured coefficients. It is seen that the first resonant peak occurs at 159 Hz coming close to the experimentally obtained value (157.2 Hz). However, the other modes are not comparatively dominant.



Figure 11. Identified bearing parameters.



Figure 12. Frequency spectra at left bearing.

CONCLUSION

An optimisation-based bearing parameter identification procedure was illustrated from the dynamic response in a dual disk rotor dynamic model. Flexible rotor system was analysed with finite element model and the responses in frequency domain have been obtained at different speeds using nonlinear bearing forces. The error function defined as a difference of actual and ideal response amplitudes was minimised with modified particle swarm optimisation scheme. The approach was further tested with vibration response data obtained from an experimental test rotor supported on hydrodynamic radial bearings. The methodology is reliable and predicts the coefficients with limited computation effort. Unlike existing methods of identification, present approach can estimate the equivalent bearing parameters from nonlinear force models and with experimental data. The average time can be further minimised by using approximate solution technique for nonlinear dynamic equations that predicts the frequency spectra without need of time responses or by using a surrogate model to avoid the time-domain simulations. It is also possible to predict cross coupling stiffness and damping coefficients as well as nonlinear stiffness coefficients at the bearings.

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APPENDIX-A

The component forces on the journal and bearing are expressed as follows:

$$\begin{cases} f_{ix} \\ f_{iy} \end{cases} = -\frac{\sqrt{(y_i + 2\dot{x}_i) + (x_i - 2\dot{y}_i)}}{(1 - x_i \cos \varphi - y_i \sin \varphi)} \begin{cases} 3x_i V_i - \sin \alpha_i G_i - 2 \cos \alpha_i F_i \\ 3y_i V_i - \cos \alpha_i G_i - 2 \sin \alpha_i F_i \end{cases}$$
(A.1)

where V, G, F, α are the lubricant force variants.

$$V_{i} = \frac{2 + (y_{i} \cos \alpha_{i} - x_{i} \sin \alpha_{i})G_{i}}{(1 - x_{i}^{2} - y_{i}^{2})}$$
(A.2)

$$G_{i} = \frac{\pi}{\sqrt{1 - x_{i}^{2} - y_{i}^{2}}} - \frac{2}{\sqrt{1 - x_{i}^{2} - y_{i}^{2}}} \tan^{-1} \left(\frac{y_{i} \cos \alpha_{i} - x_{i} \sin \alpha_{i}}{\sqrt{1 - x_{i}^{2} - y_{i}^{2}}} \right)$$
(A.3)

$$F_{i} = \frac{(x_{i} \cos \alpha_{i} + y_{i} \sin \alpha_{i})}{(1 - x_{i}^{2} - y_{i}^{2})}$$
(A.4)

$$\alpha_{i} = \tan^{-1} \left(\frac{y_{i} + 2\dot{x}_{i}}{x_{i} - \dot{y}_{i}} \right) - \frac{\pi}{2} \operatorname{sign} \left(\frac{y_{i} + 2\dot{x}_{i}}{x_{i} - y_{i}} \right) - \frac{\pi}{2} \operatorname{sign}(y_{i} + 2\dot{x}_{i})$$
(A.5)

$$\begin{cases} f_{ox} \\ f_{oy} \end{cases} = -\frac{\sqrt{(y_R + 2\dot{x}_R) + (x_R - 2\dot{y}_R)}}{(1 - x_R \cos \varphi - y_R \sin \varphi)} \begin{cases} 3x_R V_0 - \sin \alpha_0 G_0 - 2 \cos \alpha_0 F_0 \\ 3y_R V_0 - \cos \alpha_0 G_0 - 2 \sin \alpha_0 F_0 \end{cases}$$
(A.6)

$$V_{o} = \frac{2 + (y_{R} \cos \alpha_{o} - x_{R} \sin \alpha_{o})G_{o}}{(1 - x_{R}^{2} - y_{R}^{2})}$$
(A.7)

$$G_{o} = \frac{\pi}{\sqrt{1 - x_{R}^{2} - y_{R}^{2}}} - \frac{2}{\sqrt{1 - x_{R}^{2} - y_{R}^{2}}} \tan^{-1} \left(\frac{y_{R} \cos \alpha_{o} - x_{R} \sin \alpha_{o}}{\sqrt{1 - x_{R}^{2} - y_{R}^{2}}} \right)$$
(A.8)

$$F_{o} = \frac{(x_{R} \cos \alpha_{o} + y_{R} \sin \alpha_{o})}{(1 - x_{R}^{2} - y_{R}^{2})}$$
(A.9)

$$\alpha_{o} = \tan^{-1}\left(\frac{y_{R}+2\dot{x}_{R}}{x_{R}-\dot{y}_{R}}\right) - \frac{\pi}{2}\operatorname{sign}\left(\frac{y_{R}+2\dot{x}_{R}}{x_{R}-y_{R}}\right) - \frac{\pi}{2}\operatorname{sign}(y_{R}+2\dot{x}_{R})$$
(A.10)